

On a Sufficient Condition for Observability of Nonlinear Switched Systems and Observer Design Strategy

Hyungbo Shim and Aneel Tanwani

Abstract—This paper presents a sufficient condition for observability of switched systems that involve state jumps and comprise nonlinear dynamical subsystems affine in control. Without assuming observability of individual modes, the sufficient condition is based on gathering partial information from each mode so that the state is recovered completely after some time. Using this result, an observer is designed which employs a novel ‘back-and-forth’ technique to generate state estimates. Under the assumption of persistent switching, analysis shows that the estimate converges asymptotically to the actual state of the system.

I. INTRODUCTION

We study observability conditions and an observer design for a class of switched nonlinear systems Σ , described as

$$\dot{x}(t) = f_{\sigma(t)}(x(t)) + g_{\sigma(t)}(x(t))u(t), \quad t \neq \{t_q\}, \quad (1a)$$

$$x(t_q) = p_{\sigma(t_q^-)}(x(t_q^-)), \quad (1b)$$

$$y(t) = h_{\sigma(t)}(x(t)), \quad (1c)$$

where $x : \mathbb{R} \mapsto \mathbb{R}^n$ is the state trajectory, $y : \mathbb{R} \mapsto \mathbb{R}^{d_y}$ is the output, the measurable function $u : \mathbb{R} \mapsto \mathbb{R}^{d_u}$ is the input belonging to some input class \mathcal{U} of interest, and $\sigma : \mathbb{R} \mapsto \mathbb{N}$ is the switching signal that is right-continuous and changes its value at switching times $\{t_q\}$, $q \in \mathbb{N}$. Let t_0 be the initial time and the jump map (1b) applies at $t = t_q$, $q \geq 1$. It is assumed that there are a finite number of switching times in any finite time interval. The switching mode σ and the switching times $\{t_q\}$ may come from a supervisory logic controller, or may be determined internally depending on the system state. In any case, we treat them as a known, external input in this paper. It is assumed that the solution $x(t)$ remains in a compact set $\mathcal{X} \subset \mathbb{R}^n$ on the time interval of interest. All the vector fields and functions are assumed to be smooth, and thus, the existence and uniqueness of the solution, for all times, are guaranteed by the fact that the solution remains in a compact set.

When dealing with observability of nonlinear systems, there are different notions that are involved. The work in [8], [10] talk about observability in local neighborhoods of the state space. Authors in [9] describe the notion of ‘large-time’ versus ‘small-time’ observability where the difference lies in whether it is possible to recover the state instantaneously in time or the system becomes observable after certain time

interval. If the system description has exogenous inputs acting on it, then the question arises whether observability holds for all inputs or not [5], [6]; if it does, the system is called uniformly observable.

The concept of observability studied in this paper is a refinement of the ‘large-time observability’ already considered in the literature (e.g., [9]) and the ‘uniform observability’ studied in [5], [6]. Switched systems can be thought of as a family of dynamical subsystems, where a switching signal determines the active subsystem at each time instant. It is entirely possible that none of these subsystems is observable in the sense that information about the full state is not immediate in the output signal [8], [10]. But the information available from each mode can be combined in a certain manner so that under some conditions, it is possible to recover the state vector completely after some time. This explains how the concept of ‘large-time’ comes into picture when dealing with switched systems and our goal is to derive conditions that make the system large-time observable on a given set \mathcal{X} . Moreover, since we are interested in an observer construction at the end, the observability for all inputs (i.e., uniform observability) is of concern in order for the observer to be independent of particular inputs.

For switched systems, among other structural properties, observability and observer design for linear case have been actively studied during the past decade. Some initial observer results on switched systems, such as [1], [13] for linear case and [3] for nonlinear case, have assumed that each mode in the system is in fact observable admitting a state observer, and have treated the switching as a source of perturbation effect. More relaxed approaches do not assume observability of the individual modes, and the notion of gaining observability by switching for linear systems has appeared in, e.g., [4], [20], [21]. Even though limited to the linear case, it is not clear how the conditions in [4], [20], [21] can lead to feasible observer design. On the other hand, there is not much literature on the observability of switched nonlinear systems.

The main contribution of this paper lies in the unified treatment of observability conditions and observer design which has not been discussed in literature for nonlinear systems, to the authors’ knowledge. For the observer design, our approach shares the same spirit as [2], and the result of this paper can be regarded as an extension of [2], in the sense that, a coordinate-independent condition is derived for observability and nonlinear systems are treated with a new observer design strategy.

The notation used in this paper is summarized as follows.

H. Shim is with ASRI, School of Electrical Engineering, Seoul National University, Korea, and is supported by the National Research Foundation of Korea grant funded by the Korea government (MEST) (grant 2011-0001276).

A. Tanwani is with Coordinated Science Laboratory, University of Illinois at Urbana-Champaign, USA, and is sponsored by NSF grant ECCS-0821153.

$\mathcal{R}(A)$ implies the range space of the columns of matrix A , and A^\top is the transpose of A . We denote $[x_1^\top, x_2^\top]^\top$ simply by $\text{col}(x_1, x_2)$, and $\lambda_{1\sim k} := \text{col}(\lambda_1, \dots, \lambda_k)$. For a signal $x(t)$, $x_{[t_1, t_2]}$ means $\{(x(t), t) : t_1 \leq t \leq t_2\}$. The differential of a map p acting on the vector field v is denoted by p_*v . For a distribution \mathcal{W} , $p_*\mathcal{W} = \{p_*v \mid v \in \mathcal{W}\}$. We call a distribution \mathcal{W} at x° ‘nonsingular’ when $\dim \mathcal{W}$ is constant in a neighborhood of x° . The notation $\text{l.com} \{\lambda_1(x), \dots, \lambda_k(x)\}$ means a set of linear combinations of the functions λ_i with constant coefficients, i.e., $\{\sum_{i=1}^k c_i \lambda_i(x) : c_i \in \mathbb{R}\}$. Now let \mathcal{X} be a set in \mathbb{R}^n , and whenever we say a property holds ‘on \mathcal{X} ,’ we mean that it holds for every $x \in \mathcal{X}$. Smooth functions $\lambda_1(x), \dots, \lambda_k(x)$, defined on \mathcal{X} , are said to be *independent on \mathcal{X}* if their differential one-forms, $d\lambda_1(x), \dots, d\lambda_k(x)$ are linearly independent on \mathcal{X} . In addition, if there exist $n-k$ smooth functions $\lambda_{k+1}, \dots, \lambda_n$ such that $\text{col}(\lambda_1(x), \dots, \lambda_n(x))$ becomes a diffeomorphism from \mathcal{X} to \mathbb{R}^n , then we say that $\lambda_1, \dots, \lambda_k$ are *potential coordinate functions on \mathcal{X}* . We also recall that the Lie derivative of a function λ along the vector field f is $L_f \lambda(x) := \frac{\partial \lambda(x)}{\partial x} f(x)$ and $L_f(d\lambda) = dL_f \lambda$. Some proofs have been omitted in this conference paper due to space constraints.

II. PRELIMINARIES

Let us formalize the notion of observability considered in this paper.

Definition 1: A system Σ with a switching signal $\sigma(\cdot)$ is *large-time uniformly observable on a set $\mathcal{X} \subset \mathbb{R}^n$* if there exist a finite time $T > t_0$ so that $x(T)$ is determined uniquely from $y_{[t_0, T]}$, $u_{[t_0, T]}$, and $\sigma_{[t_0, T]}$ for any measurable input $u_{[t_0, T]}$, when the state $x(t)$ remains in \mathcal{X} for $t \in [t_0, T]$. If the time $T > t_0$ can be chosen arbitrarily, then the system Σ is called *small-time uniformly observable on a set \mathcal{X}* . \triangleleft

In case of no jump map (1b), the knowledge of $x(T)$, $\sigma_{[t_0, T]}$, and $u_{[t_0, T]}$ determines $x_{[t_0, T]}$ uniquely. This is not the case in general because the jump map (1b) may not be reversible. From the definition, if a certain mode of system Σ is small-time observable and the switching signal activates that mode at a certain time, then the system is automatically large-time observable. Note that $x(T)$ may be reconstructed using the derivatives of $y(\cdot)$ and $u(\cdot)$ (although differentiation should not be used in the observer construction). It is noted that, although the observability in Definition 1 is uniform with respect to the input u , uniformity with respect to the switching signal σ is not required.

The following lemma will be frequently used in the paper.

Lemma 1: Consider a codistribution \mathcal{W} generated by exact one-forms, i.e., $\mathcal{W} = \text{span} \{d\lambda_1, \dots, d\lambda_k\}$ with $1 \leq k \leq n$, where $\lambda_1, \dots, \lambda_k$ are smooth potential coordinate functions defined on a set $\mathcal{X} \subset \mathbb{R}^n$.

- 1) If the codistribution \mathcal{W} is invariant with respect to a smooth vector field $f(x)$, i.e.,

$$L_f \mathcal{W} \subset \mathcal{W}$$

on \mathcal{X} , then there exists a smooth vector field $F : \lambda_{1\sim k}(\mathcal{X}) \rightarrow \mathbb{R}^k$ such that $\frac{\partial \lambda_{1\sim k}}{\partial x} \Big|_x \cdot f(x) = F(\lambda_{1\sim k}(x))$ on \mathcal{X} .

- 2) If a smooth function $h : \mathcal{X} \rightarrow \mathbb{R}$ satisfies

$$dh \in \mathcal{W}$$

on \mathcal{X} , then, there exists a smooth function $H : \lambda_{1\sim k}(\mathcal{X}) \rightarrow \mathbb{R}$ such that $h(x) = H(\lambda_{1\sim k}(x))$ on \mathcal{X} .

- 3) Let \mathcal{W}' be another codistribution such that $\dim(\mathcal{W} + \mathcal{W}')$ is constant on \mathcal{X} , and suppose that $\mathcal{W} + \mathcal{W}' = \text{span} \{d\lambda_{1\sim k}, d\mu'_j : j = 1, \dots, \bar{r}'\}$ where $\bar{r}' = \dim(\mathcal{W} + \mathcal{W}') - \dim \mathcal{W}$, and the elements of $\{\lambda_{1\sim k}, \mu'_{1\sim \bar{r}'}\}$ are smooth potential coordinate functions on \mathcal{X} . If a smooth map $p : \mathcal{X} \rightarrow \mathbb{R}^n$ satisfies

$$p_*(\ker \mathcal{W} \cap \ker \mathcal{W}') \subset \ker \mathcal{W}$$

on \mathcal{X} , then there exists a smooth map $P : \lambda_{1\sim k}(\mathcal{X}) \times \mu'_{1\sim \bar{r}'}(\mathcal{X}) \rightarrow \mathbb{R}^k$ such that $\lambda_{1\sim k}(p(x)) = P(\lambda_{1\sim k}(x), \mu'_{1\sim \bar{r}'}(x))$ while x and $p(x)$ are contained in \mathcal{X} . \triangleleft

Before dealing with the switched case, let us consider the system (1a) and (1c) for a fixed mode q , without the jump map (1b) for now. In particular, we note that the individual subsystems may not be observable, which calls for the classical Kalman decomposition [10]: changing the coordinates so that the system is explicitly split into the observable part and the unobservable part. For this, let the observation space \mathcal{O}_q be the linear space of functions over \mathbb{R} containing all $h_{q,i}$'s (where $h_{q,i}$ is the i -th element of h_q) and all repeated Lie derivatives $L_{v_1} L_{v_2} \dots L_{v_k} h_{q,i}$ with $v_l \in \{f, g_{q,1}, \dots, g_{q,d_u}\}$ ($g_{q,j}$ is the j -th column of g_q).

Assumption 1: For each mode q , the codistribution $d\mathcal{O}_q = \text{span} \{d\lambda : \lambda \in \mathcal{O}_q\}$ has constant dimension k_q ; $\dim d\mathcal{O}_q = k_q$ on the set \mathcal{X} . In addition, there are k_q smooth potential coordinate functions $\lambda_{q,j}$, $j = 1, \dots, k_q$, such that

$$d\mathcal{O}_q = \text{span} \{d\lambda_{q,1}, \dots, d\lambda_{q,k_q}\} \quad \text{on } \mathcal{X}. \quad \triangleleft$$

Under Assumption 1, additional functions $\lambda_{q,k_q+1}, \dots, \lambda_{q,n}$ can be found to yield a diffeomorphism $\lambda_q := \text{col}(\lambda_{q,1}(x), \dots, \lambda_{q,n}(x))$ on \mathcal{X} . Then, Lemma 1.1 implies that the system (1a) and (1c) for mode q can be written in the new coordinates $\text{col}(\xi'_q, \xi_q)$ with $\xi'_q = \lambda_{q,(k_q+1)\sim n}(x) \in \mathbb{R}^{n-k_q}$ and $\xi_q = \lambda_{q,1\sim k_q}(x) \in \mathbb{R}^{k_q}$ as

$$\dot{\xi}'_q = F'_q(\xi'_q, \xi_q) + G'_q(\xi'_q, \xi_q)u \quad (2a)$$

$$\dot{\xi}_q = F_q(\xi_q) + G_q(\xi_q)u \quad (2b)$$

$$y = H_q(\xi_q). \quad (2c)$$

This representation is valid on the set $\lambda_q(\mathcal{X})$.

Assumption 2: The reduced-order subsystem (2b) and (2c) is small-time uniformly observable on the set $\Xi_q := \lambda_{q,1\sim k_q}(\mathcal{X})$, which is the projection of $\lambda_q(\mathcal{X})$ onto the ξ_q -coordinates. \triangleleft

Remark 1: As a matter of fact, the simple condition that $\dim d\mathcal{O}_q(x) = k_q$ in a neighborhood of some $x^\circ \in \mathcal{X}$ guarantees, by Frobenius theorem, that there exists a local neighborhood $\mathcal{X}'(\subset \mathcal{X})$ of x° such that Assumption 1 holds. Compared to this local observability (studied in, e.g., [8], [10], [12]), Assumptions 1 and 2 may be thought of as global versions (“global” in the sense of the whole region \mathcal{X}). \triangleleft

As long as we restrict our attention to (2b) and (2c) for each mode, Assumption 2 becomes the standard uniform observability assumption that has often been studied in the literature (see [7] and references therein). Assumption 2 can be checked in various ways; for instance, if the class of inputs \mathcal{U} consists of smooth functions only, then one may try to find a function \mathcal{E} such that

$$\xi_q = \mathcal{E}(y, \dot{y}, \dots, y^{(n_y-1)}, u, \dot{u}, \dots, u^{(n_u-1)})$$

where $n_y \in \mathbb{N}$ and $n_u \in \mathbb{N}$, and that the function $\mathcal{E}(\cdot, u, \dot{u}, \dots, u^{(n_u-1)})$ is surjective onto Ξ_q for all $u(\cdot) \in \mathcal{U}$. The existence of such a function \mathcal{E} to reconstruct $x(T)$ is used as the definition of uniform observability in [10], [19]. Other ways to check Assumption 2 can be found in [14].

III. SUFFICIENT CONDITION FOR OBSERVABILITY

In deriving the sufficient condition for observability, we do not assume the individual modes of the system to be observable. Hence, in order to recover the system state $x(t)$, partial information obtained from each mode is accumulated. This partial information is quantified in terms of the maximal integral submanifold of the distribution $d\mathcal{O}_q^\perp$ which has the property that the states on the slices (or “leaves”) of this submanifold are not distinguishable by the output of mode q . As soon as a switch occurs, the indistinguishable states must be contained in the intersection of the integral submanifolds of both modes before and after the switching. In this way, switching reduces the uncertainty in the state. Continuing in this manner, with subsequent switching, we expect to reduce the size of submanifold that contains the indistinguishable states. Eventually, if the corresponding intersections reduce to a point, we obtain observability. However, while the intersections are taken *at the same time*, the information contained in the integral submanifolds is scattered in time because each one of them becomes available sequentially as time goes on. This suggests that the partial information, obtained at each mode, should evolve uncorrupted along the dynamics of subsequent modes until all the information is gathered to compute the state. Inspired by this intuition, we present structural conditions which guarantee that the evolution of the partial information is feasible without being affected by the unknown quantities in subsequent modes.

Before presenting the condition, let us rename the switching sequence for convenience. That is, when the switching signal $\sigma(t)$ takes the mode sequence $\{q_1, q_2, \dots\}$, we rename them as increasing integers $\{1, 2, 3, \dots\}$ which is ever increasing even though the same mode is revisited. This description also incorporates cases where there is a state jump without change in dynamics or the mode change does not involve state jumps.

Theorem 1: Suppose that Assumptions 1 and 2 hold, and define $d\mathcal{O}'_q := \text{span} \{d(\lambda_{q,i} \circ p_{q-1}) : i = 1, \dots, k_q\}$ for each $q \geq 2$. On \mathcal{X} , define a sequence of codistributions \mathcal{W}_q , with $\mathcal{W}_0 := \{0\}$, as:

\mathcal{W}_q is the largest nonsingular and involutive codistribution, invariant with respect to f_q and g_q , con-

tained in $(d\mathcal{O}_q + \mathcal{W}_{q-1})$ such that $(p_q)_*(\ker \mathcal{W}_q \cap \ker d\mathcal{O}'_{q+1}) \subset \ker \mathcal{W}_q$.

If

- 1) $\exists m \geq 1$ such that, on \mathcal{X} ,

$$\dim(d\mathcal{O}_m + \mathcal{W}_{m-1}) = n,$$

(or simply $\dim \mathcal{W}_m = n$ because $\mathcal{W}_m = d\mathcal{O}_m + \mathcal{W}_{m-1}$ by construction),

- 2) the codistributions \mathcal{W}_q ($1 \leq q \leq m$), $d\mathcal{O}_q + \mathcal{W}_{q-1}$ ($2 \leq q \leq m$), and $\mathcal{W}_q + d\mathcal{O}'_{q+1}$ ($1 \leq q \leq m-1$) are nonsingular on \mathcal{X} . Moreover,

- (a) \exists smooth potential coordinate functions $\{\phi_{q,i}, \omega_{q,j} : i = 1, \dots, \bar{k}_q, j = 1, \dots, \bar{l}_q, \bar{k}_q + \bar{l}_q = \dim \mathcal{W}_q\}$ on \mathcal{X} such that

$$\begin{aligned} \mathcal{W}_q = \text{span} \{d\phi_{q,1}, \dots, d\phi_{q,\bar{k}_q}, d\omega_{q,1}, \dots, d\omega_{q,\bar{l}_q}\}, \\ d\phi_{q,i} \in d\mathcal{O}_q, \quad d\omega_{q,j} \notin d\mathcal{O}_q, \end{aligned} \quad (3)$$

- (b) \exists smooth potential coordinate functions $\{\mu_{q,i} : i = 1, \dots, \bar{r}_q, \bar{r}_q = \dim(d\mathcal{O}_q + \mathcal{W}_{q-1})\}$ on \mathcal{X} such that

$$d\mathcal{O}_q + \mathcal{W}_{q-1} = \text{span} \{d\mu_{q,1}, \dots, d\mu_{q,\bar{r}_q}\}, \quad (4)$$

$$\begin{aligned} \mu_{q,i} \in \text{l.com} \{\lambda_{q,1}, \dots, \lambda_{q,k_q}, \phi_{q-1,1}, \dots, \\ \phi_{q-1,\bar{k}_{q-1}}, \omega_{q-1,1}, \dots, \omega_{q-1,\bar{l}_{q-1}}\}, \end{aligned} \quad (5)$$

- (c) \exists smooth potential coordinate functions $\{\mu'_{q,j} : j = 1, \dots, \bar{r}'_q, \bar{r}'_q = \dim(\mathcal{W}_q + d\mathcal{O}'_{q+1}) - \dim \mathcal{W}_q\}$ on \mathcal{X} such that

$$\begin{aligned} \mathcal{W}_q + d\mathcal{O}'_{q+1} = \text{span} \{d\phi_{q,1}, \dots, d\phi_{q,\bar{k}_q}, \\ d\omega_{q,1}, \dots, d\omega_{q,\bar{l}_q}, d\mu'_{q,1}, \dots, d\mu'_{q,\bar{r}'_q}\}, \end{aligned} \quad (6)$$

$$\mu'_{q,j} \in \text{l.com} \{\lambda_{q+1,1} \circ p_q, \dots, \lambda_{q+1,k_{q+1}} \circ p_q\}, \quad (7)$$

then the system (1) is large-time uniformly observable on \mathcal{X} for all the switching signals containing the consecutive subsequence $\{1, 2, \dots, m\}$. \triangleleft

Following observations are immediate: (a) $d\mathcal{O}_q$ itself is invariant with respect to f_q and g_q by construction, (b) if $p_q(x) = x$, so that there is no state jump, then the condition $(p_q)_*(\ker \mathcal{W}_q \cap \ker d\mathcal{O}'_{q+1}) \subset \ker \mathcal{W}_q$ automatically holds, (c) the “largest” codistribution in the assumption of Theorem 1 is well-defined, because involutivity and invariance of a codistribution generated by exact one-forms is preserved under the addition, and if two smooth nonsingular codistributions \mathcal{W}_a and \mathcal{W}_b satisfy $p_*(\ker \mathcal{W}_i \cap \mathcal{D}) \subset \ker \mathcal{W}_i$, where $i \in \{a, b\}$, for any differentiable map p and any distribution \mathcal{D} , then $p_*(\ker(\mathcal{W}_a + \mathcal{W}_b) \cap \mathcal{D}) \subset \ker(\mathcal{W}_a + \mathcal{W}_b)$.

The compactness of the set \mathcal{X} guarantees the solution without finite escape time, and will be used for observer construction in the next section. If all the assumptions hold with $\mathcal{X} = \mathbb{R}^n$, then the observability property becomes global in case the solution has no finite escape time. On the other hand, if local observability is of interest, then the assumptions get simpler by removing the items 2(a), 2(b), and 2(c).

Corollary 1: Suppose that Assumptions 1 and 2 hold in a neighborhood of a point $x^o \in \mathcal{X}$. If each of the codistributions \mathcal{W}_q , $d\mathcal{O}_q + \mathcal{W}_{q-1}$, and $\mathcal{W}_q + d\mathcal{O}'_{q+1}$ are

nonsingular at x^o , \mathcal{W}_q is smooth and involutive at x^o , and $\dim(d\mathcal{O}_m + \mathcal{W}_{m-1})(x^o) = n$, then the system is large-time uniformly observable in some neighborhood of x^o for all the switching signals containing the consecutive subsequence $\{1, 2, \dots, m\}$. \triangleleft

Now we present the proof of Theorem 1, which is constructive in the sense that a technique to recover $x(t)$ at some time $t = T > t_{m-1}$ is revealed (rather than discussing the indistinguishability of two different states). This way paves the road to the observer design in the next section.

Proof: [Proof of Theorem 1] Consider the interval prior to the first switching, $[t_0, t_1)$. Since $\mathcal{W}_1 \subset d\mathcal{O}_1$, we have that

$$\mathcal{W}_1 = \text{span} \{d\phi_{1,1}, d\phi_{1,2}, \dots, d\phi_{1,\bar{k}_1}\}, \quad \bar{k}_1 \leq k_1.$$

Because $d\phi_{1,i}$, for each $i = 1, \dots, \bar{k}_1$, is an element of $d\mathcal{O}_1$ that is generated by the differentials of $\lambda_{1,l}$, $l = 1, \dots, k_1$, and $\lambda_{1,l}$'s are potential coordinate functions on \mathcal{X} (by Assumption 1), $\phi_{1,i}$ is a function of $\lambda_{1,1 \sim k_1}$ only (by Lemma 1.2). Since $\xi_1 := \lambda_{1,1 \sim k_1}(x)$ is small-time uniformly observable on $\lambda_{1,1 \sim k_1}(\mathcal{X})$ (by Assumption 2), the value of the vector $\xi_1(t) = \lambda_{1,1 \sim k_1}(x(t))$, and thus, $\phi_{1,1 \sim \bar{k}_1}(x(t))$ are recovered for $t \in [t_0, t_1)$.

Now Lemma 1.3, with the item 2(c) and $(p_1)_*(\ker \mathcal{W}_1 \cap \ker d\mathcal{O}_2) \subset \ker \mathcal{W}_1$, implies the existence of a function \check{P}_1 (and then P_1 below, since $\phi_{1,i}$ is a function of $\lambda_{1,1 \sim k_1}$) such that

$$\begin{aligned} \phi_{1,1 \sim \bar{k}_1}(x(t_1)) &= \phi_{1,1 \sim \bar{k}_1}(p_1(x(t_1^-))) \\ &= \check{P}_1(\phi_{1,1 \sim \bar{k}_1}(x(t_1^-)), \mu'_{1,1 \sim \bar{r}'_1}(x(t_1^-))) \\ &= P_1(\lambda_{1,1 \sim k_1}(x(t_1^-)), \lambda_{2,1 \sim k_2} \circ p_1(x(t_1^-))) \\ &= P_1(\lambda_{1,1 \sim k_1}(x(t_1^-)), \lambda_{2,1 \sim k_2}(x(t_1))), \end{aligned} \quad (8)$$

where the third equality follows from (3) and (7). Next, consider the interval $[t_1, t_2)$. For $i = 1, \dots, \bar{k}_2$, using Lemma 1.2, the condition $d\phi_{2,i} \in d\mathcal{O}_2 = \text{span} \{d\lambda_{2,1}, \dots, d\lambda_{2,k_2}\}$ guarantees that $\phi_{2,i}$ is a function of $\lambda_{2,1 \sim k_2}$ only. Again by Assumption 2, the vector $\xi_2(t) := \lambda_{2,1 \sim k_2}(x(t))$, and thus, $\phi_{2,1 \sim \bar{k}_2}(x(t))$ are recovered for the interval $[t_1, t_2)$.

Now observing that \mathcal{W}_2 is invariant w.r.t. f_2 and g_2 , $\mathcal{W}_2 = \text{span} \{d\phi_{2,1}, \dots, d\phi_{2,\bar{k}_2}, d\omega_{2,1}, \dots, d\omega_{2,\bar{l}_2} : \bar{k}_2 + \bar{l}_2 = \dim \mathcal{W}_2\}$, and $\{\phi_{2,1 \sim \bar{k}_2}, \omega_{2,1 \sim \bar{l}_2}\}$ are potential coordinate functions on \mathcal{X} , we apply Lemma 1.1 and obtain smooth vector fields F_2^* and G_2^* such that, with $z_2 := \omega_{2,1 \sim \bar{l}_2}(x)$,

$$\begin{aligned} \dot{z}_2 &= \frac{\partial \omega_{2,1 \sim \bar{l}_2}}{\partial x}(x) \cdot (f(x) + g(x)u) \\ &= \check{F}_2^*(z_2, \phi_{2,1 \sim \bar{k}_2}(x)) + \check{G}_2^*(z_2, \phi_{2,1 \sim \bar{k}_2}(x))u \\ &= F_2^*(z_2, \lambda_{2,1 \sim k_2}(x)) + G_2^*(z_2, \lambda_{2,1 \sim k_2}(x))u \\ &= F_2^*(z_2, \xi_2) + G_2^*(z_2, \xi_2)u, \end{aligned} \quad (9)$$

over $[t_1, t_2)$. In this interval, the vector $\xi_2(t) = \lambda_{2,1 \sim k_2}(x(t))$ is recovered. Hence, if the initial condition $z_2(t_1) = \omega_{2,1 \sim \bar{l}_2}(x(t_1))$ is recovered, the vector $z_2(t)$ on the interval $[t_1, t_2)$ is also known by solving the differential equation (9).

Note that $d\omega_{2,j} \in \mathcal{W}_2 \subset (d\mathcal{O}_2 + \mathcal{W}_1) = \text{span} \{d\mu_{2,1}, \dots, d\mu_{2,\bar{r}_2}\}$, $j = 1, \dots, \bar{l}_2$, by the definition

of \mathcal{W}_2 and the item 2(b). Therefore, by Lemma 1.2, $\omega_{2,j}$ can be written as a function of $\mu_{2,i}$'s, which leads to

$$\begin{aligned} z_2(t_1) &= \omega_{2,1 \sim \bar{l}_2}(x(t_1)) = \check{S}_2^*(\mu_{2,1 \sim \bar{r}_2}(x(t_1))) \\ &= S_2^*(\xi_2(t_1), \phi_{1,1 \sim \bar{k}_1}(x(t_1))) \\ &= S_2^*(\xi_2(t_1), P_1(\lambda_{1,1 \sim k_1}(x(t_1^-)), \lambda_{2,1 \sim k_2}(x(t_1)))) \\ &= S_2^*(\xi_2(t_1), P_1(\xi_1(t_1^-), \xi_2(t_1))) \end{aligned} \quad (10)$$

in which, the third equality uses (5) with ω_1 being null, and the fourth equality follows from (8).

This process is repeated to find F_q^* , G_q^* , S_q^* , and P_q . For instance, we can find P_2 such that

$$\begin{aligned} \begin{bmatrix} \phi_{2,1 \sim \bar{k}_2}(x(t_2)) \\ \omega_{2,1 \sim \bar{l}_2}(x(t_2)) \end{bmatrix} &= \begin{bmatrix} \phi_{2,1 \sim \bar{k}_2}(p_2(x(t_2^-))) \\ \omega_{2,1 \sim \bar{l}_2}(p_2(x(t_2^-))) \end{bmatrix} \\ &= \check{P}_2(\phi_{2,1 \sim \bar{k}_2}(x(t_2^-)), \omega_{2,1 \sim \bar{l}_2}(x(t_2^-)), \mu'_{2,1 \sim \bar{r}'_2}(x(t_2^-))) \\ &= P_2(\lambda_{2,1 \sim k_2}(x(t_2^-)), \omega_{2,1 \sim \bar{l}_2}(x(t_2^-)), \lambda_{3,1 \sim k_3} \circ p_2(x(t_2^-))) \\ &= P_2(\xi_2(t_2^-), z_2(t_2^-), \xi_3(t_2)), \end{aligned}$$

and find S_3^* such that

$$\begin{aligned} z_3(t_2) &= \omega_{3,1 \sim \bar{l}_3}(x(t_2)) = \check{S}_3^*(\mu_{3,1 \sim \bar{r}_3}(x(t_2))) \\ &= S_3^*(\xi_3(t_2), \phi_{2,1 \sim \bar{k}_2}(x(t_2)), \omega_{2,1 \sim \bar{l}_2}(x(t_2))) \\ &= S_3^*(\xi_3(t_2), P_2(\xi_2(t_2^-), z_2(t_2^-), \xi_3(t_2))). \end{aligned}$$

In summary, for each time interval $[t_{q-1}, t_q)$, $q = 1, \dots, m$, we have the differential equation (with z_1 being null)

$$\dot{\xi}_q = F_q(\xi_q) + G_q(\xi_q)u, \quad \xi_q \in \mathbb{R}^{k_q}, \quad (11a)$$

$$y = H_q(\xi_q), \quad (11b)$$

$$\dot{z}_q = F_q^*(z_q, \xi_q) + G_q^*(z_q, \xi_q)u, \quad z_q \in \mathbb{R}^{\bar{l}_q}, \quad (11c)$$

$$z_q(t_{q-1}) = S_q^*(\xi_q(t_{q-1}), P_{q-1}(\xi_{q-1}(t_{q-1}^-), z_{q-1}(t_{q-1}^-), \xi_q(t_{q-1}))), \quad (11d)$$

in which, $\xi_q(t)$ and $z_q(t)$ are completely determined.

At any time $t = T > t_{m-1}$, it follows under the assumption in item 1, i.e. $\dim \mathcal{W}_m = n$, that the vectors $\xi_m(T)$ and $z_m(T)$ are completely recovered or equivalently $\text{col}(\phi_{m,1 \sim \bar{k}_m}(x(T)), \omega_{m,1 \sim \bar{l}_m}(x(T)))$ is determined. In this way $x(T)$ is recovered uniquely by the inverse, because $\text{col}(\phi_{m,1 \sim \bar{k}_m}, \omega_{m,1 \sim \bar{l}_m})$ is a diffeomorphism. \blacksquare

IV. OBSERVER DESIGN

Based on the study of large-time observability, let us now discuss the design of an asymptotic observer for the system (1). By asymptotic observer, we mean that the estimate $\hat{x}(t)$ that it generates, converges to the plant state $x(t)$ as time tends to infinity. In order to achieve this, we introduce the following assumptions.

Assumption 3: 1) The switching is persistent and happens within the duration D ; that is,

$$t_q - t_{q-1} \leq D, \quad \forall q \in \mathbb{N} \quad (12)$$

where t_q is the switching time.

- 2) The solution $x(t)$ of the plant (1) remains in a compact set $\mathcal{X} \in \mathbb{R}^n$, and the input $u(t)$ is uniformly bounded; $|u(t)| \leq M_u$.
- 3) There is an $m \in \mathbb{N}$ such that the assumption of Theorem 1 holds on a set $\bar{\mathcal{X}}$ that properly contains \mathcal{X} , and the mode sequence repeats the particular modes $\{1, 2, \dots, m\}$; that is, $\sigma(t) = ((q-1) \bmod m) + 1$ for $[t_{q-1}, t_q)$, $q \in \mathbb{N}$. \triangleleft

We do not consider the time consumed for computation by assuming that the data processor is fairly fast compared to the plant process. The computation time, however, needs to be considered in real-time application if the plant itself is fast.

The observer we propose is of hybrid-type, and has the form

$$\begin{aligned} \dot{\hat{x}}(t) &= \bar{f}_q(\hat{x}(t)) + \bar{g}_q(\hat{x}(t))u(t), & t \in [t_{q-1}, t_q), \\ \hat{x}(t_q) &= \begin{cases} \bar{p}_q(\hat{x}(t_q^-)), & (q \bmod m) \neq 0, \\ \bar{p}_q(\mathcal{L}_q(y_{[t_{q-m}, t_q)}, u_{[t_{q-m}, t_q)})), & (q \bmod m) = 0, \end{cases} \end{aligned} \quad (13)$$

with an initial condition $\hat{x}(t_0) \in \mathcal{X} \subset \bar{\mathcal{X}}$, where \bar{f}_q , \bar{g}_q , and \bar{p}_q are globally Lipschitz and they have the same values as f_q , g_q , and p_q , respectively, inside the compact set \mathcal{X} . Their global Lipschitz property can always be obtained by modifying them outside the set \mathcal{X} , using the so-called ‘Lipschitz extension’.¹ It is seen that the observer consists of a plant copy with an estimate update law by some operator \mathcal{L}_q , which we design in this section. In fact, we present a design of \mathcal{L}_q , using some dynamic observers for partial states at each mode and an inversion algorithm logic in order to achieve,

$$|\tilde{x}(t_m)| \leq \gamma |\tilde{x}(t_0)| \quad (14)$$

where $0 < \gamma < 1$ and $\tilde{x} := \hat{x} - x$. The Lipschitz property of (13) and Assumptions 3.1 and 3.2 guarantee that $\sup_{t \in [t_{(j-1)m}, t_{jm})} |\tilde{x}(t)| \leq \Gamma |\tilde{x}(t_{(j-1)m})|$ with a constant Γ and $j \in \mathbb{N}$. In this way, if (14) holds then its repeated application leads to $\lim_{t \rightarrow \infty} |\tilde{x}(t)| = 0$.

The proposed observer construction is based on the representation (11) of the plant (1). The idea is that, for each interval $[t_{q-1}, t_q)$, $q = 1, \dots, m$, a conventional nonlinear observer, which we call ξ_q -observer, is employed to obtain the estimate $\hat{\xi}_q(t)$ for that interval. At the same time, a z_q -observer, replicating (11c) and (11d), is constructed as follows:

$$\dot{\hat{z}}_q = \bar{F}_q^*(\hat{z}_q, \hat{\xi}_q) + \bar{G}_q^*(\hat{z}_q, \hat{\xi}_q)u, \quad 2 \leq q \leq m, \quad (15)$$

with the initial condition given by:

$$\hat{z}_q(t_{q-1}) = \bar{S}_q^*(\hat{\xi}_q(t_{q-1}), \bar{P}_{q-1}(\hat{\xi}_{q-1}(t_{q-1}^-), \hat{\xi}_q(t_{q-1}))), \quad (16)$$

and $\hat{z}_1 := 0$ for convenience. Here \bar{F}_q^* is the Lipschitz extension of F_q^* with respect to the set $Z_q \times \Xi_q$ and so

¹Since the plant state $x(t)$ remains in \mathcal{X} , this modification can also be applied to the plant model (1). See [5] for its utility in observer construction. Detailed procedures to the modification have been discussed in [14], [15].

on (Z_q is the image of \mathcal{X} through $\omega_{q,1 \sim \bar{l}_q}$, and Ξ_q through $\lambda_{q,1 \sim k_q}$). In fact, the variable z_q is not an observable quantity for the mode q . Intuitively speaking, the role of z_q -observer is not to reduce the error $\hat{z}_q(t) := \hat{z}_q(t) - z_q(t)$, but to deliver the estimates $\hat{\xi}_{q-1}(t_{q-1}^-)$ and $\hat{z}_{q-1}(t_{q-1}^-)$, that are obtained from the previously active mode and are encoded in the initial condition (16), to the next mode through $\hat{z}_q(t)$ along the system dynamics. Suppose that, seen at time $t = t_m^-$, an ideal observer provides the exact information of $\xi_q(t)$ on each interval $[t_{q-1}, t_q)$, $q = 1, \dots, m$, using the stored input u and the output y with the model (11a) and (11b). For example, with exact values of $\hat{\xi}_1(t_1^-) = \xi_1(t_1^-)$ and $\hat{\xi}_2(t_1) = \xi_2(t_1)$, we obtain the exact value of $\hat{z}_2(t_1) = z_2(t_1)$ by (16). Then, integration of (15) for $q = 2$ results in exact values of $\hat{z}_2(t) = z_2(t)$ on $[t_1, t_2)$. This process repeats until we get $\hat{\xi}_m(t_m^-) = \xi_m(t_m^-)$ and $\hat{z}_m(t_m^-) = z_m(t_m^-)$. With Assumption 3.3, i.e. $\dim \mathcal{W}_m = n$, $x(t_m^-)$ is now determined uniquely from $\xi_m(t_m^-)$ and $z_m(t_m^-)$, as the map

$$x(t_m^-) \mapsto \begin{bmatrix} \phi_{m,1 \sim \bar{k}_m}(x(t_m^-)) \\ \omega_{m,1 \sim \bar{l}_m}(x(t_m^-)) \end{bmatrix} = \begin{bmatrix} \chi(\xi_m(t_m^-)) \\ z_m(t_m^-) \end{bmatrix} \quad (17)$$

is invertible; here χ is a function such that $\chi(\lambda_{m,1 \sim k_m}(x)) = \phi_{m,1 \sim \bar{k}_m}(x)$ whose existence is guaranteed by Lemma 1.2. For convenience let us denote the inverse map by Ψ , so that $x(t_m^-) = \Psi(\xi_m(t_m^-), z_m(t_m^-))$. As a result, we choose the estimate update law in (13) to be:

$$\hat{x}(t_m^-) = \bar{\Psi}(\hat{\xi}_m(t_m^-), \hat{z}_m(t_m^-)) =: \mathcal{L}_q(y_{[t_0, t_m)}, u_{[t_0, t_m)}) \quad (18)$$

where $\bar{\Psi}$ is Lipschitz extension of Ψ . Through this relation, the plant state is recovered as $\hat{x}(t_m^-) = x(t_m^-)$ with exact information $\hat{\xi}_m(t_m^-) = \xi_m(t_m^-)$ and $\hat{z}_m(t_m^-) = z_m(t_m^-)$.

However, asymptotic observers in practice inevitably introduce some error in $\hat{\xi}_q(t)$ while estimating $\xi_q(t)$. Moreover, the estimation of $\xi_q(t)$ on the *entire interval* $[t_{q-1}, t_q)$ needs more attention because the conventional observers, initiated at the time $t = t_{q-1}$, often experience the transient overshoot before they converge to the proper estimates. Reducing the transient period by increasing observer gain may worsen the situation because of the peaking phenomenon [17]; that is, the peaking in $\hat{\xi}_q(t)$ may damage the role of (15) because large error in $|\hat{z}_q(t_q^-) - z_q(t_q^-)|$ may occur in spite of small error in $|\hat{z}_q(t_{q-1}) - z_q(t_{q-1})|$.

We overcome this obstacle by employing a new back-and-forth estimation technique. Suppose that the ξ_q -observer over the interval $[t_{q-1}, t_q)$ is written as,

$$\begin{aligned} \dot{\hat{\xi}}_q^f &= \hat{F}_q(\hat{\xi}_q^f) + \hat{G}_q(\hat{\xi}_q^f)u + K_q^f(\hat{\xi}_q^f, u, y)(y - \hat{H}_q(\hat{\xi}_q^f)), \\ \hat{\xi}_q^f(t_{q-1}) &= \bar{\lambda}_{q,1 \sim k_q}(\hat{x}(t_{q-1})), & 1 \leq q \leq m \end{aligned} \quad (19)$$

where $\bar{\lambda}_{q,1 \sim k_q}$ is the Lipschitz extension of $\lambda_{q,1 \sim k_q}$ and the superscript ‘ f ’ indicates ‘forward’, whose meaning will soon become clear from the context. Here, \hat{F}_q , \hat{G}_q , and \hat{H}_q can be any modified vector fields or functions as long as $\hat{F}_q(\xi_q) = F_q(\xi_q)$, $\hat{G}_q(\xi_q) = G_q(\xi_q)$, and $\hat{H}_q(\xi_q) = H_q(\xi_q)$ for $\xi_q \in \Xi_q$. In this way, we allow many observer design techniques

available in the literature (which may utilize the Lipschitz extension of F_q , G_q , and H_q in another coordinate system).

Assumption 4: Under the assumption that $\xi_q(t) \in \Xi_q$ for $[t_{q-1}, t_q]$, $q = 1, \dots, m$, the subsystem (11a) and (11b) admits an observer of the form (19), which can be made to converge to the state $\xi_q(t)$ arbitrarily fast, that is, for arbitrarily small constants $b > 0$ and $c > 0$, there exist an injection gain $K_q^f(\cdot)$ and a class- \mathcal{KL} function $\beta^{f,q}$ satisfying

$$\begin{aligned} \beta^{f,q}(a, t) &< ca \quad \text{for all } a > 0 \text{ and } b \leq t \leq \tau_q \quad (20) \\ |\hat{\xi}_q^f(t) - \xi_q(t)| &\leq \beta^{f,q}(|\hat{\xi}_q^f(t_{q-1}) - \xi_q(t_{q-1})|, t - t_{q-1}) \quad (21) \end{aligned}$$

for $t \in [t_{q-1}, t_q]$. \triangleleft

Remark 2: Many results in the literature, such as [5], [11], yield an observer satisfying Assumption 4 with $\beta^{f,q}$ being an exponential function. \triangleleft

Now consider another (backward) observer described as

$$\begin{aligned} \dot{\xi}_q^b &= -\hat{F}_q(\xi_q^b) - \hat{G}_q(\xi_q^b)u(t_q - t) \\ &\quad - K_q^b(\xi_q^b, u(t_q - t), y(t_q - t))(y(t_q - t) - \hat{H}_q(\xi_q^b)), \\ \xi_q^b(0) &= \hat{\xi}_q^f(t_q^-), \quad t \in (0, \tau_q], \quad 1 \leq q \leq m. \quad (22) \end{aligned}$$

Actually, the trajectory $\xi_q^b(t) := \xi_q(t_q - t)$ satisfies the differential equation

$$\dot{\xi}_q^b = -F_q(\xi_q^b) - G_q(\xi_q^b)u(t_q - t), \quad y(t_q - t) = H_q(\xi_q^b),$$

with $\xi_q^b(0) = \xi_q(t_q^-)$, for $t \in (0, \tau_q]$, and therefore, (22) can be thought of as one possible observer for it. We further assume that Assumption 4 holds for this case as well, with $\hat{\xi}_q^f$, ξ_q , $\beta^{f,q}$ and K_q^f replaced by $\hat{\xi}_q^b$, ξ_q^b , $\beta^{b,q}$ and K_q^b , respectively. Once Assumption 4 holds for (19), this additional requirement is mild. For example, the designs of [11] and [5] readily satisfy this requirement. Then, using the input u and the output y stored over the interval $[t_{q-1}, t_q]$, we run the observer (19) first from the initial condition $\hat{\xi}_q^f(t_{q-1}) = \bar{\lambda}_{q,1 \sim k_q}(\hat{x}(t_{q-1}))$, followed by integrating (22) from 0 to τ_q . After that, we take our final estimate $\hat{\xi}_q(t)$ as

$$\hat{\xi}_q(t) = \begin{cases} \hat{\xi}_q^b(t_q - t), & t \in [t_{q-1}, t_{q-1} + \tau_q/2), \\ \hat{\xi}_q^f(t), & t \in [t_{q-1} + \tau_q/2, t_q). \end{cases} \quad (23)$$

From Assumption 4 let us assume that, with $b = \tau_q/2$ and a given $c \in (0, 1)$, both K_q^f and K_q^b are designed. With $\tilde{\xi}_q := \hat{\xi}_q - \xi_q$, $\tilde{\xi}_q^f := \hat{\xi}_q^f - \xi_q$, and $\tilde{\xi}_q^b := \hat{\xi}_q^b - \xi_q^b$, it is seen that

$$\begin{aligned} \sup_{t \in [t_{q-1} + \frac{\tau_q}{2}, t_q)} |\tilde{\xi}_q(t)| &= \sup_{t \in [t_{q-1} + \frac{\tau_q}{2}, t_q)} |\tilde{\xi}_q^f(t)| \\ &\leq \sup_{t \in [t_{q-1} + \frac{\tau_q}{2}, t_q)} \beta^{f,q}(|\tilde{\xi}_q^f(t_{q-1})|, t - t_{q-1}) \leq c|\tilde{\xi}_q^f(t_{q-1})|, \\ \sup_{t \in [t_{q-1}, t_{q-1} + \frac{\tau_q}{2})} |\tilde{\xi}_q(t)| &= \sup_{t \in [t_{q-1}, t_{q-1} + \frac{\tau_q}{2})} |\tilde{\xi}_q^b(t_q - t)| \\ &\leq \sup_{t \in [t_{q-1}, t_{q-1} + \frac{\tau_q}{2})} \beta^{b,q}(\beta^{f,q}(|\tilde{\xi}_q^f(t_{q-1})|, \tau_q), t_q - t) \\ &\leq c^2|\tilde{\xi}_q^f(t_{q-1})| \leq c|\tilde{\xi}_q^f(t_{q-1})|. \end{aligned}$$

Therefore, implementation of the observers in (19) and (22) leads to

$$\begin{aligned} \sup_{t \in [t_{q-1}, t_q)} |\tilde{\xi}_q(t)| &\leq c|\tilde{\xi}_q^f(t_{q-1})| \\ &= c|\bar{\lambda}_{q,1 \sim k_q}(\hat{x}(t_{q-1})) - \lambda_{q,1 \sim k_q}(x(t_{q-1}))|. \quad (24) \end{aligned}$$

Theorem 2: Under the assumptions so far, there is $c^* > 0$ such that, for each $c \in (0, c^*)$, the inequality (14) holds. \triangleleft

This theorem concludes the asymptotic error convergence to zero. (Proof is omitted but available from the authors.)

ACKNOWLEDGEMENT

The authors are grateful to Daniel Liberzon for helpful discussions on this paper.

REFERENCES

- [1] A. Alessandri and P. Coletta. Switching observers for continuous-time and discrete-time linear systems. In *American Control Conference*, pages 2516–2521, 2001.
- [2] A. Balluchi, L. Benvenuti, M. D. Di Benedetto, and A. L. Sangiovanni-Vincentelli. Observability for hybrid systems. In *IEEE Conf. on Decision and Control*, pages 1159–1164, 2003.
- [3] J. P. Barbot, H. Saadaoui, M. Djemai, and N. Manamanni. Nonlinear observer for autonomous switching systems with jumps. *Nonlinear Analysis: Hybrid Systems*, 1:537–547, 2007.
- [4] P. Collins and J. H. van Schuppen. Observability of piecewise-affine hybrid systems. In R. Alur and G. J. Pappas, editors, *Hyb. Sys: Comp. & Control*, Vol. 2993 of *LNCS*, pages 265–279. Springer, Berlin, 2004.
- [5] J. P. Gauthier, H. Hammouri, and S. Othman. A simple observer for nonlinear systems. *IEEE Trans. on Automatic Control*, 37:875–880, 1992.
- [6] J. P. Gauthier and I. Kupka. Observability and observers for nonlinear systems. *SIAM J. on Control*, 32(4):975–994, 1994.
- [7] J. P. Gauthier and I. Kupka. *Deterministic observation theory and applications*. Cambridge University Press, 2001.
- [8] R. Hermann and A. Krener. Nonlinear controllability and observability. *IEEE Trans. on Automatic Control*, 22(5):728–740, 1977.
- [9] J. P. Hespanha, D. Liberzon, D. Angeli, and E. D. Sontag. Nonlinear norm-observability notions and stability of switched systems. *IEEE Trans. on Automatic Control*, 50(2):154–168, 2005.
- [10] A. Isidori. *Nonlinear Control System*. Springer, 3rd edition, 1995.
- [11] A. J. Krener and A. Isidori. Linearization by output injection and nonlinear observers. *Systems & Control Letters*, 3:47–52, 1983.
- [12] H. Nijmeijer and A. J. van der Schaft. *Nonlinear Dynamical Control Systems*. Springer, 1990.
- [13] S. Pettersson. Designing switched observers for switched systems using multiple Lyapunov functions and dwell-time switching. In *IFAC Conf. on Analysis and Design of Hybrid Systems*, pages 18–23, 2006.
- [14] H. Shim. *A passivity-based nonlinear observer and a semi-global separation principle*. PhD thesis, Seoul National University, 2000.
- [15] H. Shim and J. H. Seo. Nonlinear output feedback stabilization on a bounded region of attraction. *Int. J. of Control*, 73(5):416–426, 2000.
- [16] Z. Sun and S. S. Ge. *Switched Linear Systems: Control and Design*. Springer-Verlag London, 2005.
- [17] H. Sussmann and P. V. Kokotovic. The peaking phenomenon and the global stabilization of nonlinear systems. *IEEE Trans. on Automatic Control*, 36(4):424–440, 1991.
- [18] A. Tanwani, H. Shim, and D. Liberzon. Observability implies observer design for switched linear systems. *Conf. on Hyb. Sys: Comp. & Control*, 2011.
- [19] A. Teel and L. Praly. Global stabilizability and observability imply semiglobal stabilizability by output feedback. *Systems & Control Letters*, 22:313–325, 1994.
- [20] R. Vidal, A. Chiuso, S. Soatto, and S. Sastry. Observability of linear hybrid systems. In O. Maler and A. Pnueli, editors, *Hyb. Sys: Comp. & Control*, Vol. 2623 of *LNCS*, pages 526–539. Springer, Berlin, 2003.
- [21] G. Xie and L. Wang. Necessary and sufficient conditions for controllability and observability of switched impulsive control systems. *IEEE Trans. on Automatic Control*, 49:960–966, 2004.