

# Planning agile motions for quadrotors in constrained environments

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**Abstract**—Planning physically realistic and easily controllable motions of flying robots requires considering dynamics. This paper presents a local trajectory planner, based on a simplified dynamic model of quadrotors, which fits the requirements to be integrated into a global motion planning approach. It relies on a closed-form solution to compute curves in the kinodynamic state space that tend to minimize the flying time. These curves have suitable continuity properties and guarantee respect of physical limits of the system (i.e. bounds for the time-derivatives of the pose coordinates). The paper explains how this local planner can be used within different motion planning approaches that enable the treatment of difficult problems in constrained environments.

## I. INTRODUCTION

When planning motions for Unmanned Aerial Vehicles (UAVs) such as quadrotors, it is important to consider the dynamic model of the system since not every geometrically valid path corresponds to a feasible motion. For example, because of its dynamic behavior, flying upside down, even for a relatively short period of time, is hardly manageable for a fixed-pitch quadrotor. Aiming to avoid the difficulties and the high computational cost involving kinodynamic motion planning [1], the problem is usually treated in two stages. The first stage applies a basic path planning approach, disregarding dynamic considerations. For this, sampling-based path planning algorithms [2], such as the Rapidly-exploring Random Tree (RRT) or the Probabilistic Roadmap (PRM), can be used to produce a collision free path for the center of mass of the robot. Indeed, since the robot orientation depends on dynamic aspects, collisions cannot be tested for the robot itself but for its smallest bounding-sphere. In a second stage, this path, usually consisting of a sequence of straight-line segments in  $\mathbb{R}^3$ , has to be transformed into a dynamic trajectory. Trajectory generation methods, such as [3], [4], [5], can be applied in this stage to each portion of the path. The overall trajectory then consists of a sequence of movements from one hovering position to another, which leads to severe sub-optimality in terms of execution time. However, such an unsuitable trajectory can be subsequently optimized using several types of algorithms.

The aforementioned *decoupled approach* is computationally efficient, and can be successfully applied to solve many motion planning problems for UAVs (see for example [6], [7], [8]). However, several classes of problems cannot be treated using this approach, because the robot orientation

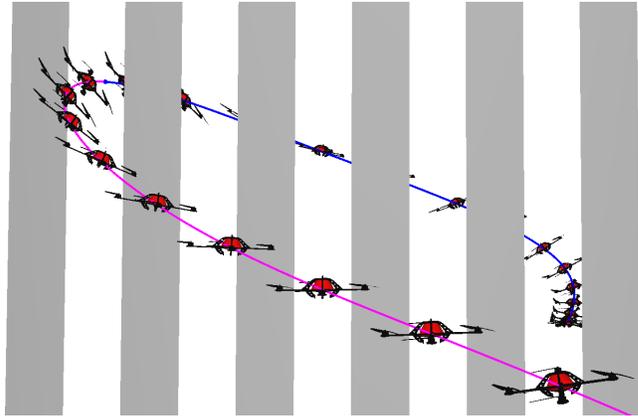


Fig. 1. Solution to the *slot problem* obtained with a basic RRT algorithm using the proposed local planner as a steering method.

cannot be properly considered at the geometric stage. One of such problems, which we refer to as the *slot problem*, is illustrated in Fig. 1. In this problem, the robot has to go through a narrow slot-shaped passage whose width is smaller than the diameter of the smallest bounding-sphere of the robot. Since a collision-free path does not exist for this sphere, the decoupled approach will fail to find a solution. Another type of problem that cannot be treated using a decoupled approach is the transportation of a large object (at medium or high velocity). When carrying such an object, it is not suitable anymore to consider the minimum bounding-sphere of the system because it does not fit its actual shape, thus leading to invalidity of far too many possible solutions. In these situations, collisions have to be tested for the actual shape of the system, and hence, non-hovering states have to be sampled and linked by collision-free flyable trajectories. According to [4], such trajectories must be smooth in the space of the flat outputs of the system with bounds on their derivatives. Therefore, motion planning in this context requires an efficient trajectory generation method able to interpolate two kinodynamic states (position, velocity and acceleration). Unfortunately, available trajectory generation methods are either computationally expensive, or unlikely for the present application, since they do not guarantee the respect of limits in velocity and acceleration, or because they require flying time as an input parameter [9].

This paper proposes a local motion (trajectory) planner that considers a simplified dynamic model of the UAV. It uses fourth-order splines to generate flyable trajectories of minimal time with respect to the proposed closed form solution. It can be used as part of an optimization method in a decoupled approach in order to generate high velocity

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trajectories in cluttered environments, or applied directly as a steering method within a sampling-based motion planner. Next sections describe the method and present results of its application to several motion planning problems in constrained environments, including the *slot problem*. These problems are also illustrated by the accompanying video.

## II. LOCAL PLANNER

### A. Overview

Thanks to the differential flatness of the quadrotor dynamics [4], any smooth trajectory in the space of flat outputs can be followed provided that derivatives are correctly bounded. Note that real physical limits of a quadrotor are its total thrust and angular accelerations (which are related to linear acceleration and snap respectively). However, as a simplification, we directly consider limitations on derivatives of the flat outputs. We also choose to limit linear velocity for safety reasons in prospect of future experimental validation.

This section proposes a local motion planner for quadrotors that, given a set of bounds for the derivatives of the flat outputs, produces such trajectories while trying to minimize the flying time. First, the dynamic model is presented together with some basic concepts. The closed-form of a specific fourth-order spline used to solve the problem independently for each output is then proposed. Finally, the way outputs are synchronized to form the final trajectory is explained.

### B. Simplified dynamic model of a quadrotor, configuration space, geometric and kinodynamic state spaces

A quadrotor can be modeled as a free-flying object in  $\mathbb{R}^3$ . It has six degrees of freedom: three for its position ( $\mathbf{OG} = \mathbf{r} \in \mathbb{R}^3$ ) and three for its orientation (pitch, roll and yaw angles:  $[\theta, \phi, \psi]$ ). Thus, its configuration space is  $C = \mathbb{R}^3 \times SO(3)$ . As illustrated in Fig. 2, a quadrotor of mass  $m$  is assumed to be submitted to three forces applied to its center of mass  $G$ : his weight  $\mathbf{W} = -mg.\mathbf{e}_z$ , a fluid friction force  $\mathbf{f} = -f.\mathbf{v}$  with  $f$  a fluid friction coefficient and  $\mathbf{v} = \dot{\mathbf{r}}$  (linear velocity of  $G$ ), and a thrust force  $\mathbf{T}$ , which is the resulting force of the action of the four rotors. The direction of the thrust vector is directly related to pitch and yaw angles. Furthermore, Newton's second law of motion shows that  $\mathbf{T}$  is a function of  $\mathbf{v}$  and  $\mathbf{a} = \ddot{\mathbf{r}}$  (linear acceleration of  $G$ ):

$$m.\mathbf{a} = -mg.\mathbf{e}_z - f.\mathbf{v} + \mathbf{T} \quad (1)$$

Let  $\mathbf{x}^G = [\mathbf{r}, \psi] \in \mathbb{R}^3 \times SO(2) = \mathcal{X}_G$  be the geometric state of the system (the flat outputs) and  $\mathbf{x}^K = [\mathbf{x}^G, \dot{\mathbf{x}}^G, \ddot{\mathbf{x}}^G] \in \mathcal{X}_G^3 = \mathcal{X}_K$  be the kinodynamic state,  $\mathcal{X}_G$  and  $\mathcal{X}_K$  being the the geometric and the kinodynamic state spaces, respectively. Using (1), attitude angles  $[\theta, \phi]$  can be computed from  $\mathbf{v}$  and  $\mathbf{a}$ . Hence, intermediate configurations  $q \in C$  can be extracted from a trajectory in  $\mathcal{X}_K$  with arbitrary resolution for collision deception.

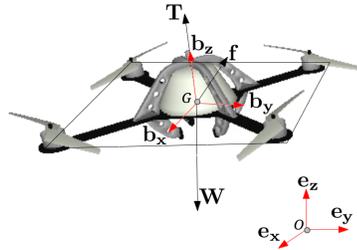


Fig. 2. Forces acting on the quadrotor in our simplified dynamic model.

### C. Closed-form solution of the trajectory generation problem

This subsection addresses the generation of smooth trajectories in the kinodynamic state space with inequality constraints on the norm of the derivatives. A trajectory between two states  $\mathbf{x}_0^K, \mathbf{x}_F^K \in \mathcal{X}_K$  is a time parametrized function  $S : [0, t_F] \rightarrow \mathcal{X}_G$  with  $t_F \in \mathbb{R}^+$  such that:

$$\begin{cases} [S(0), \dot{S}(0), \ddot{S}(0)] = \mathbf{x}_0^K \in \mathcal{X}_K \\ [S(t_F), \dot{S}(t_F), \ddot{S}(t_F)] = \mathbf{x}_F^K \in \mathcal{X}_K \end{cases} \quad (2)$$

Because the trajectory has to be continuous in  $C$ , and in order to guarantee continuity of the angular velocity,  $S$  has to be of class  $C^3$ :  $S \in C^3([0, t_F], \mathcal{X}_G)$ . This is due to the relationship between the attitude and the linear acceleration imposed by (1). In addition, since these (local) trajectories will be the elements of the solution provided by the (global) motion planner, we want to ensure continuity of the jerk (first derivative of acceleration) between two consecutive trajectories. Thus, another constraint has to be imposed:

$$\ddot{\ddot{S}}(0) = \ddot{\ddot{S}}(t_F) = 0 \quad (3)$$

The trajectory generation problem can be independently solved for each component  $i = 1 \dots 4$  of the geometric state (i.e. each flat output). The parameter  $t_{F,i} \in \mathbb{R}^+$  and  $S_i \in C^3([0, t_{F,i}], \mathbb{R})$  have to be computed for each component, taking into account bounds on the derivatives. These bounds are directly considered up to the third order. Nevertheless, since jerk has to be continuous and bounded, snap (second derivative of acceleration) has to be bounded too. In summary, the set of constraints is:

$$\begin{cases} S_i(0) = x_0, \dot{S}_i(0) = v_0 \\ \dot{S}_i(0) = a_0, \ddot{S}_i(0) = 0 \\ S_i(t_{F,i}) = x_F, \dot{S}_i(t_{F,i}) = v_F \\ \dot{S}_i(t_{F,i}) = a_F, \ddot{S}_i(t_{F,i}) = 0 \\ \forall t \in [0, t_{F,i}], \begin{cases} |\dot{S}_i(t)| \leq v_{max} \in \mathbb{R}^{*+} \\ |\ddot{S}_i(t)| \leq a_{max} \in \mathbb{R}^{*+} \\ |\ddot{\ddot{S}}_i(t)| \leq j_{max} \in \mathbb{R}^{*+} \\ |\ddot{\ddot{S}}_i(t)| \leq s_{max} \in \mathbb{R}^{*+} \end{cases} \end{cases} \quad (4)$$

In order to minimize  $t_{F,i}$ , full speed has to be reached as soon as possible and maintained as long as possible. Let  $D$  refer to this phase of constant velocity and let  $v_D \in [-v_{max}, v_{max}]$  be the value of velocity during this phase. This means that time spent to reach it has to be minimized: velocity have to get from  $v_0$  to  $v_D$  and then from  $v_D$  to  $v_F$  as quickly as possible. Hence time spent at maximum acceleration during those two phases has to be maximized. It

implies to minimize the durations of acceleration variations. Let  $B$  refer to the phase of constant acceleration during the first velocity variation and  $a_B \in [-a_{max}, a_{max}]$  be this constant acceleration. Let  $G$  be the phase of constant acceleration during the second velocity variation and  $a_G \in [-a_{max}, a_{max}]$  this constant acceleration. The phase in which acceleration varies from  $a_0$  to  $a_B$  is noted  $A$ , and  $C$  is used to refer to the phase in which acceleration varies from  $a_B$  to zero. The phases where acceleration varies from zero to  $a_G$  and from  $a_G$  to  $a_F$  are noted  $E$  and  $H$  respectively. All these phases are illustrated with an example in Fig. 3. The principle holds for higher derivatives: time spent at maximum jerk during acceleration variations has to be maximized and durations of jerk variations have to be minimized. This implies to maximize time spent at maximum snap during jerk variations and to minimize the durations of snap variations. This very last part is easy since, following our approach, snap can be discontinuous: snap variation is actually a snap commutation of duration zero ( $\forall t \in [0, t_{F,i}], \ddot{S}_i(t) \in \{-s_{max}, 0, s_{max}\}$ ). In other words,  $\ddot{S}_i$  is a piecewise constant function, which implies that  $S_i$  is a piecewise polynomial function (a spline) of the fourth order.

In the discussion above, we have mentioned seven main phases in the (local) trajectory. For three of them ( $B, D, G$ ), acceleration is zero and hence snap is also zero. The four others ( $A, C, E, H$ ) are divided into three sub-phases (1, 2, 3). Phases 1 and 3 correspond to jerk variations and then  $|\ddot{S}_i(t)| = s_{max}$ . In phase 2, jerk is constant hence snap is zero. The sign of the snap is opposed during phases 1 and 3, which have the same duration. Tab. I presents the notation used for the durations of the phases and the expressions of the snap as functions of  $a_B$  and  $a_G$ .

The expression of the snap during phase  $A_1$  is explained as follows: if  $a_B > a_0$ ,  $\ddot{S}_i$  has to be increasing, which implies that  $\ddot{S}_i$  has to be positive. Since  $\ddot{S}_i(0) = 0$ ,  $\ddot{S}_i$  has to be increasing during phase  $A_1$ , and thus  $\ddot{S}_i(t) = +s_{max}$ . Following the same reasoning,  $\ddot{S}_i(t) = -s_{max}$  if  $a_B < a_0$ . If  $a_B = a_0$ , then phase  $A$  is not needed, and thus  $\ddot{S}_i(t) = 0$ . A similar reasoning can be applied to understand the expressions for phases  $C, E$  and  $H$ .

Note that this closed-form solution is an intuitive way of minimize flying time. However, there is no guarantee of optimality because of our choices for simplification. The actual optimal solution could be found by applying Pontryagin maximum principle to the corresponding optimal control problem (snap as scalar control input and inequality constraints on the state) but it would imply heavy numerical computation, making it not suited to our needs. The sub-optimality factor induced by our choices is yet to be measured. From now on, both *optimal time/velocity* are to be read as: *optimal time/velocity with respect to our closed-form solution*.

#### D. Duration of the phases

At this point, the local trajectory for one coordinate (i.e its corresponding spline) is defined by  $a_B, a_G$  and the duration

TABLE I  
VALUE OF THE SNAP AND DURATION OF EACH PHASE OF THE SPLINE

Phase	Value of the Snap	Duration
$A_1$	$\text{sign}(a_B - a_0) \cdot s_{max}$	$t_{A,1}$
$A_2$	0	$t_{A,2}$
$A_3$	$\text{sign}(a_0 - a_B) \cdot s_{max}$	$t_{A,1}$
$B$	0	$t_B$
$C_1$	$-\text{sign}(a_B) \cdot s_{max}$	$t_{C,1}$
$C_2$	0	$t_{C,2}$
$C_3$	$\text{sign}(a_B) \cdot s_{max}$	$t_{C,1}$
$D$	0	$t_D$
$E_1$	$\text{sign}(a_G) \cdot s_{max}$	$t_{E,1}$
$E_2$	0	$t_{E,2}$
$E_3$	$-\text{sign}(a_G) \cdot s_{max}$	$t_{E,1}$
$G$	0	$t_G$
$H_1$	$\text{sign}(a_F - a_G) \cdot s_{max}$	$t_{H,1}$
$H_2$	0	$t_{H,2}$
$H_3$	$\text{sign}(a_G - a_F) \cdot s_{max}$	$t_{H,1}$

of all phases. This subsection explains how it can be defined using a single parameter:  $v_D$ . First,  $t_{A,1}, t_{A,2}, t_{C,1}$  and  $t_{C,2}$  can be expressed as functions of  $a_B, a_0, j_{max}$  and  $s_{max}$ . The same goes for  $t_{H,1}, t_{H,2}, t_{E,1}$  and  $t_{E,2}$  as functions of  $a_G, a_F, j_{max}$  and  $s_{max}$ . We provide next explanations for  $t_{A,1}$  and  $t_{A,2}$ . Let us define  $\delta_{B0} = \text{sign}(a_B - a_0)$ . From Tab. I and (4), we can write:

$$a_B = \delta_{B0} \cdot s_{max} \cdot t_{A,1}^2 + \delta_{B0} \cdot s_{max} \cdot t_{A,1} \cdot t_{A,2} + a_0 \quad (5)$$

and  $\ddot{S}_i(t_{A,1}) = \delta_{B0} \cdot s_{max} \cdot t_{A,1}$ , which implies  $|\ddot{S}_i(t_{A,1})| = s_{max} \cdot t_{A,1} \leq j_{max}$ , and thus:

$$s_{max} \cdot t_{A,1}^2 \leq \frac{j_{max}^2}{s_{max}} \quad (6)$$

Let us note  $a_{lim} = \frac{j_{max}^2}{s_{max}}$ . Phase  $A_2$  is only needed when phases  $A_1$  and  $A_3$  are not enough to reach  $a_B$ . If there is no phase  $A_2$ , then  $t_{A,2} = 0$ . Then, using (5) and (6), we can write:  $|a_B - a_0| = s_{max} \cdot t_{A,1}^2 \leq a_{lim}$ . If  $|a_B - a_0| > a_{lim}$ , then  $t_{A,2} \neq 0$  (phase  $A_2$  is needed), and  $t_{A,1} = \frac{j_{max}}{s_{max}}$  (its maximum value). In that case,  $t_{A,2} = \frac{|a_B - a_0|}{j_{max}} - \frac{j_{max}}{s_{max}}$ . If  $t_{A,2} = 0$ , then  $t_{A,1} = \sqrt{\frac{|a_B - a_0|}{s_{max}}}$ . The principle is the same for phases  $C, E$  and  $H$  and formulas are similar (details are not provided here due to limited space). For phase  $C$ ,  $|a_B - a_0|$  is replaced by  $|a_B|$ , by  $|a_G|$  in phase  $E$  and by  $|a_G - a_F|$  in phase  $H$ . The spline is now defined by  $(a_B, a_G, t_B, t_D, t_G)$ . Both couples  $(a_B, t_B)$  and  $(a_G, t_G)$  can actually be expressed as a function of  $v_D$ . Using Tab. I and (4),  $v_D$  can be expressed as a function of  $(a_B, t_{A,1}, t_{A,2}, t_B, t_{C,1}, t_{C,2})$  and hence as a function of  $(a_B, t_B)$ . Let  $V_B : (a_B, t_B) \mapsto v_D$  be this function. Let us also define  $v_B^{min} = V_B(-a_{max}, 0)$  and  $v_B^{max} = V_B(a_{max}, 0)$ . It is possible to show (not detailed here due to limited space) that, for all  $v_D$  in  $[v_B^{min}, v_B^{max}]$ , there is a unique  $a_B \in ([-a_{max}, \min(0, a_0)] \cup [\max(0, a_0), a_{max}])$  such that  $V_B(a_B, 0) = v_D$ . If  $v_D < v_B^{min}$ , then  $a_B = -a_{max}$  and phase B is needed. Its duration is simply  $t_B = |v_D - v_B^{min}|/a_{max}$ . If  $v_D > v_B^{max}$ , then  $a_B = a_{max}$  and  $t_B = |v_D - v_B^{max}|/a_{max}$ . Hence, as previously announced  $(a_B, t_B)$  can be expressed as a function of  $v_D$ . The same goes for  $(a_G, t_G)$ .

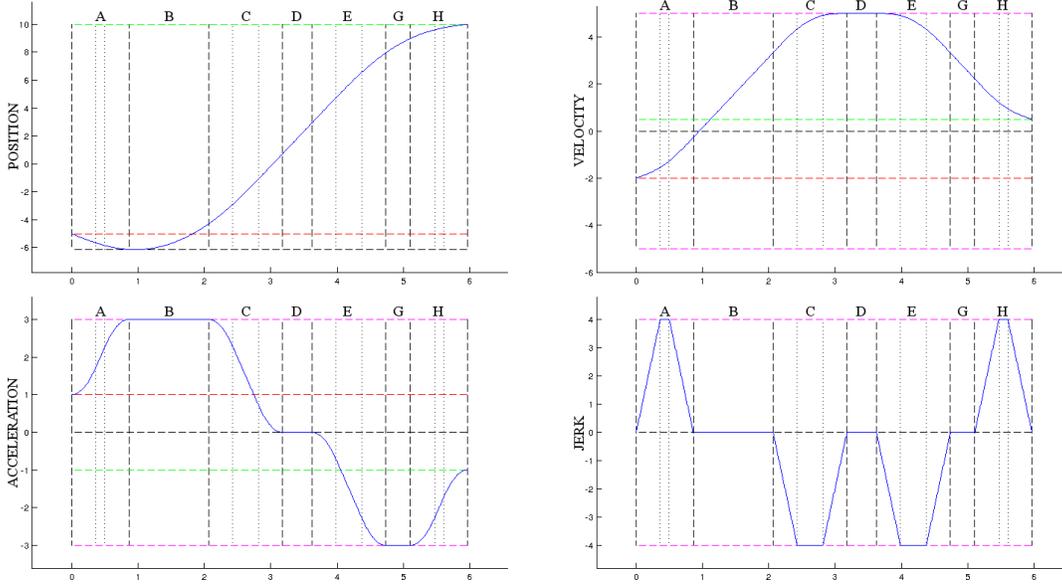


Fig. 3. Example of trajectory provided by the local planner for one coordinate. The bounds of each derivative are represented by pink dashed lines. Red and green dashed lines represent initial and final values, respectively. In phase A,  $a_B$  is reached with respect to maximum jerk and snap constraints. In phase B, the system travels towards  $v_D$  at constant acceleration  $a_B$ . C is a deceleration phase and D the phase of constant velocity. H, G and E are the symmetric phases of A, B and C, respectively. Note that, in order to show all the phases,  $|v_D| = v_{max}$  and  $|a_B| = |a_D| = a_{max}$  in this example. In a general case, some phases can have zero duration.

At this point the spline is defined by  $(v_D, t_D)$ . Let  $t_C$  be the value of the time parameter at the end of phase C, and  $t_E$  its value at the beginning of phase E. As explained previously, these two values can be expressed as a function of  $v_D$ . This is also the case for  $S_i(t_C)$  and  $S_i(t_E)$ . Let  $\Delta_S$  be the function  $v_D \mapsto S_i(t_E) - S_i(t_C)$ . If  $v_D \neq 0$ , then  $t_D = \Delta_S(v_D)/v_D$ . Therefore, the spline is completely defined by the sole value of  $v_D$ . Note that, necessarily:

$$v_D \cdot \Delta_S(v_D) \leq 0 \quad (7)$$

#### E. Optimal velocity

Once the spline is defined by  $v_D$ , this section explains how to compute the optimal value  $v_{opt}$  that minimizes  $t_{F,i}(v_D)$ . Let  $\mathcal{V}_{valid}$  be the set of values of  $v_D$  that meet the constraint (7) and let us note  $\mathcal{V}_0 = [\min(0, v_{opt}), \max(0, v_{opt})]$ . For synchronization purposes, it is necessary that  $\mathcal{V}_0 \subset \mathcal{V}_{valid}$ , which implies:

$$\forall v \in \mathcal{V}_0, v \cdot \Delta_S(0) \leq 0 \quad (8)$$

Let us note  $\delta_0 = \text{sign}(\Delta_S(0))$ . Since minimizing  $t_{F,i}$  requires maximizing  $|v_D|$ , (8) implies that, if  $\Delta_S$  has zeros between  $v = 0$  and  $v = \delta_0 \cdot v_{max}$ , then  $v_{opt}$  is the one of lowest absolute value. If not,  $v_{opt} = \delta_0 \cdot v_{max}$ .

#### F. Synchronization

The trajectory generation problem is now solved for each component independently. Hence, there are four different values  $(t_{F,i})_i, i = 1 \dots 4$ . This subsection explains how to synchronize these one-dimensional trajectories. Solving each problem provides a couple  $(t_{F,i}, v_{opt,i})$  and the associated interval  $\mathcal{V}_{0,i}$ . The purpose of the definition of  $v_{opt}$  provided in the previous subsection is to guarantee that  $v_D \mapsto t_{F,i}(v_D)$  is continuously strictly monotonic on  $\mathcal{V}_{0,i}$ . Furthermore,

$\lim_{v_D \rightarrow 0} t_{F,i}(v_D) = +\infty$ . This implies that, for any  $t$  in  $[t_{F,i}, +\infty[$ , there is a unique  $v \in \mathcal{V}_{0,i}$  such that  $t_{F,i}(v) = t$ . This property is used to synchronize the components. The slowest component (of index  $i_{slow}$ ) is identified and a simple dichotomous search is used to find the unique  $v_D \in \mathcal{V}_{0,i}$  such that  $t_{F,i}(v_D) = t_{F,i_{slow}}(v_{opt}^{i_{slow}})$  for the other components. As a result, all the components have the same final time  $t_{F,i_{slow}}$ .

### III. TWO GLOBAL PLANNERS

#### A. Decoupled approach

This subsection explains how the local planner presented in the previous section can be used in a decoupled approach to motion (trajectory) planning. A decoupled approach consists of two stages: 1) planning a geometrically valid path in  $\mathcal{X}_G$  for the minimum bounding sphere of the quadrotor; 2) transforming this path into a trajectory in  $\mathcal{X}_K$ . In our current implementation, a bi-directional RRT algorithm is applied to explore  $\mathcal{X}_G$ , and linear interpolation is used to connect sampled states. Thus, the resulting path is a concatenation of  $n$  collision-free straight line segments in  $\mathcal{X}_G$ ,  $\{x_0^G, x_F^G\}_i, i = 1 \dots n$ . The local trajectory planner is then applied to each pair of intermediate states  $\{[x_0^G, 0, 0], [x_F^G, 0, 0]\}_i$ . The resulting trajectory is collision-free, since the center of mass of the quadrotor follows the initial geometric path. However, it is far from being time-optimal because of the imposed stops at the ends of local paths. A trajectory optimization method can be applied to reduce the time of this initial solution. We have implemented a simple but efficient method based on the random shortcut algorithm [10]. This iterative, anytime algorithm works as follows: at each iteration, two states,  $x_1^K$  and  $x_2^K$ , are randomly selected from the overall trajectory. Let us call  $x_A^K$  the initial state of the local trajectory in which



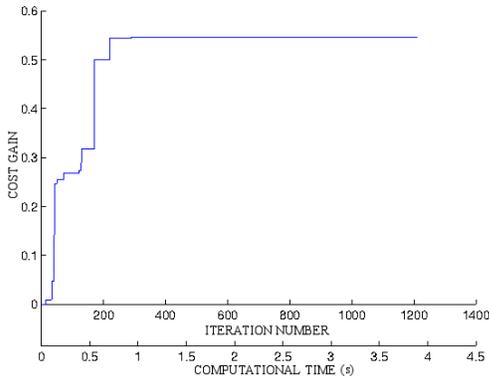


Fig. 5. Performances of the trajectory optimization method in terms of iteration number and computational time.

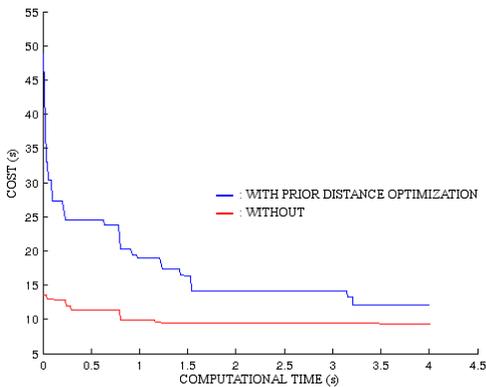


Fig. 6. Evolution of the cost during optimization with and without prior distance optimization.

illustrated in Fig. 5, about 300 iterations of the shortcut algorithm (performed in about 1 second) are sufficient for convergence. The plot represents the cost gain (i.e. the flying time decrease). Flying time of the optimized trajectory is 8.3 seconds.

An alternative strategy to obtain a time-optimized trajectory with a decoupled approach could be to optimize first the length of the geometric path. This could be done using the RRT\* algorithms instead of RRT (as proposed in [8]), or applying an optimization post-processing (i.e. path smoothing) to the solution obtained with RRT. However, results of our experiments show that this is a wrong strategy. Fig. 6 shows that the convergence of the trajectory optimization process is much slower when starting from the shortest path than when using the “rough” path provided by RRT.

### C. Direct use as a steering method: the slot problem

Finally, we present results obtained using the proposed local planner as steering method inside RRT. The method was applied to solve the difficult *slot problem* illustrated in Fig. 1. The environment is a box of  $10 \times 10 \times 3$  meters, divided in two halves by a series of aligned obstacles separated by 0.40 meters. The start position is in one side of this series of obstacles and the goal is on the other side. The diameter of the bounding sphere being of 0.54 meters, this problem can not be solved by the decoupled approach, as there is no collision-free path for the bounding sphere. Solving this problem requires taking the orientation of the quadrotor into

account, which implies considering dynamics for motion planning. Bounds for the time-derivatives of the position are the same than in the previous example. The trajectory represented in Fig. 1 has been obtained in 90 seconds with a basic RRT algorithm using our local planner as steering method. This computational time can be explained by the fact that, while going through the slot, velocity and acceleration of the quadrotor must have very specific values taken in a small interval of possibilities. The empty space between two obstacles in the configuration space becomes a very narrow passage when mapped into the kinodynamic state space.

## V. CONCLUSION

We have presented a method to compute physically-realistic local trajectories of quadrotors. The method is computationally efficient, which enables its integration into several types of global motion planning approaches. Used within a decoupled approach, based on geometric path planning and subsequent trajectory generation and optimization, motion planning problems in moderately-constrained environments can be solved in very short computing time. Very difficult problems, requiring a direct exploration of the kinodynamic state space, can also be solved in reasonably-short computing time using the proposed local planner as a steering method. As future work, we intend to further investigate the issue of finding an appropriate metric in the kinodynamic state space that could speed up the RRT algorithm. We also plan to compare the trajectories obtained with of our local planner to solutions provided by optimal control methods in order to better evaluate their quality, as it is done in [12]. Finally, experimental validation will soon be conducted in our testbed.

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